

On local existence and blow-up of solution for the higher-order nonlinear Kirchhoff-type equation with nonlinear strongly damped terms

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Abstract— In this paper ,we deal with the initial boundary value problems for higher -order kirchhoff-type equation with nonlinear strongly dissipation:

At first ,we prove the local existence and uniqueness of the solution by Galerkin method then and contracting mapping principle .Furthermore,we prove the global existence of solution , At last,we consider that blow up of solution in finite time under suitable condition .

Index Terms— Higher-order nonlinear Kirchhoff wave equation;local existence; The existence and uniqueness; blow-up

I. INTRODUCTION

In this paper we concerned with global existence and blow-up of solution for the following for Higher-order Kirchhoff-type equation with nonlinear strongly dissipation :

$$u_{tt} + (-\Delta)^m u_t + \phi(\|\nabla^m u\|^2)(-\Delta)^m u + h(u_t) = f(u) \quad (1.1)$$

$$u(x,t) = 0, \quad \frac{\partial^i u}{\partial \nu^i} = 0, \quad i = 1, 2, \dots, m-1, x \in \partial\Omega, t > 0, \quad (1.2)$$

$$u(x,0) = u_0, \quad u_t = u_1(x), \quad x \in \partial\Omega \quad (1.3)$$

Where $\Omega \subset \mathbb{R}^2$ is bounded open domain with smooth boundary; ν is the outer norm vector; $m > 1$ is a positive integer, $\phi(r)$ is a nonnegative locally Lipschitz, $h(u_t)$ is a nonlinear forcing, $f(u)$ is a nonlinear C^1 -function, $(-\Delta)^m u_t$ is a strongly dissipation.

There have been many researches on the global and blow-up solution for Kirchhoff equation. we can see [1-8].

Kosuke Ono [9] deals with the higher-order kirchhoff-type equation with nonlinear dissipation:

$$u_{tt} + M(\|\nabla u\|^2)(-\Delta)u + |u_t|^p u_t = f(u), \quad x \in \Omega, \quad t > 0, \quad (1.4)$$

$$u(x,0) = u_0(x), \quad u_t(x,0) = u_1(x), \quad \text{and} \quad u(x,t)|_{\partial\Omega} = 0 \quad (1.5)$$

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In where Ω is a bounded domain in \mathbb{R}^n with smooth boundary $\partial\Omega$, $M(r)$ is a nonnegative C^1 -function for $r \geq 0$ positive satisfying $M(\|\nabla u\|^2) = a + b\|\nabla u\|^{2\gamma}$ with $a, b \geq 0, a + b > 0$ and $\gamma \geq 1$, and $f(u)$ constants is nonlinear C^1 -function satisfying $|f(u)| \leq k_1|u|^{\alpha+1}$ and $|f'(u)| \leq k_2|u|^\alpha$ with $k_1, k_2 > 0$ and $\alpha > 0$. he investigate on global existence and blow up of solution .

Yaojun Ye [10] also studied the initial-boundary value problem for a class of nonlinear higher-order kirchhoff-type equation with dissipation term

$$u_{tt} + \left\| A^{\frac{1}{2}} u \right\|^{2p} Au + a|u_t|^{q-2} u_t = b|u|^{r-2} u, \quad x \in \Omega, t > 0, \quad (1.6)$$

$$u(x,0) = u_0(x), \quad u_t(x,0) = u_1(x), \quad x \in \Omega. \quad (1.7)$$

$$u(x,t) = 0, \quad \frac{\partial^i u(x,t)}{\partial \nu^i} = 0, \quad i = 1, 2, \dots, m-1, \quad x \in \partial\Omega, t \geq 0. \quad (1.8)$$

In a bounded domain, where $A = (-\Delta)^m, m > 1$ is a positive integer, $a, b, p > 0$ and $q, r > 2$ are constants, he obtains the global existence of solutions by construction a stable set in $H_0^m(\Omega)$ and show the energy decay estimate by applying a lemma of V.Komornik.

Takeshi Taniguchi [11] considered the existence of local solution to a weakly damped wave equation of equation of kirchhoff type the damped term and the source term

$$u_{tt}(t) - M(\|u(t)\|^2)\Delta u(t) - \gamma_2 u_t(t) + |u_t(t)|^p u_t(t) = |u(t)|^q u(t) \quad (1.9)$$

$$u(x,0) = u_0(x), \quad u_t(x,0) = u_1(x), \quad \text{and} \quad u(x,t)|_{\partial\Omega} = 0, \quad p, q, \gamma_2 > 0 \quad (1.10)$$

with an initial value $u(0) = u_0, u_t(0) = u_1$ and the Dirichlet

boundary condition $u(x,t)|_{\partial\Omega} = 0$, where Ω is an open

bounded domain in \mathbb{R}^n with smooth boundary and $M(r)$ is a locally Lipschitz function, he discuss the global existence and exponential asymptotic behavior of solution.

Recently, Gongwei Liu [12] concerned with the study of damped wave equation of kirchhoff type

$$u_{tt} - M(\|\nabla u\|_2^2)\Delta u + u_t = g(u). \quad (1.11)$$

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in $\Omega \times (0, \infty)$, with initial and Dirichlet boundary condition. where Ω is the bound of R^2 have a smooth boundary $\partial\Omega$. under the assumption that g is a function with exponential growth at infinity, he proves global existence and the decay property as well as blow-up of solutions in finite time under suitable conditions.

The paper is arranged as follows. in section 2, we state some preliminaries and some lemma ; in section 3, we obtain the existence and uniqueness of the local solution by the Banach contraction mapping principle ; in section 4, we discuss global existence and nonexistence results ; in section 5, we discuss the blow-up properties of solution for positive and negative initial energy and estimate for blow-up time T^* .

II. PRELIMINARIES

Throughout this paper ,for convenience, we denote the norm and scalar product in $L^2(\Omega)$ by $\|\cdot\|$ and (\cdot, \cdot) ; we use the next notations

$$(u, v) = \int_{\Omega} uv dx$$

$$\|u\|^2 = \int_{\Omega} |u|^2 dx$$

$$\|u\|_r = \left(\int_{\Omega} |u|^r dx \right)^{\frac{1}{r}}$$

where $C_i (i = 0 \dots 12)$, are constants,

In this section, we present some materials needed in the proof of our results, we assume that

there exists some constants

$$(G_1) \quad |f(u)| \leq k_1|u| + k_2|u|^{r+1}, \quad k_1, k_2, r > 0.$$

(G₂) there exist some constants

$$(1) \quad h(s)s \geq 0, \quad s \in R,$$

$$(2) \quad \|h(u_t)\| \leq C_0(1 + \|\nabla^m u_t\|)^{1-\sigma}, \quad 0 < \sigma < 1.$$

$$(3) \quad |h(u_t)| \leq k_3|u_t|^{p+1}, \quad k_3, p \geq 0.$$

(G₃) there exist a constant $\beta > 0$, such that

$$f(u)u \geq (2 + 4\beta)G(s) \quad \text{for all } s \in R.$$

In this paper ,we will use the next well-know Lemmas
Lemma2.1[10] Let s be a number with

$$2 \leq s < +\infty, n \leq 2m \quad \text{and} \quad 2 \leq s \leq \frac{2n}{n-2m}.$$

then there is a constant C depending on Ω and s such that

$$\|u\|_s \leq \kappa \left\| (-\Delta)^{\frac{m}{2}} u \right\|, \quad \forall u \in H_0^m(\Omega) \tag{2.1}$$

Lemma2.2[13](young inequality)Let a, b and μ are positive constants, such that

$$0 \leq r < \frac{2m}{n-2m} \quad \text{we have}$$

$$ab \leq \frac{\mu^p}{p} a^p + \frac{1}{q\mu^q} b^q \quad \left(\frac{1}{p} + \frac{1}{q} = 1, \quad p > 1, q > 1 \right) \tag{2.2}$$

Lemma2.3[14] Suppose that $\delta > 0$ and $B(t)$ is a nonnegative $C^2(0, +\infty)$ function such that

$$B''(t) - 4(\delta + 1)B'(t) + 4(\delta + 1)B(t) \geq 0. \tag{2.3}$$

if

$$B'(0) > r_2 B(0) + K_0 \tag{2.4}$$

Then we have $\forall t > 0, B'(t) > K_0$, where K_0 is a constant

$$r_2 = 2(\delta + 1) - 2\sqrt{(\delta + 1)\delta}$$

is smaller root of the equation

$$r^2 - 4(\delta + 1)r + 4(\delta + 1) = 0 \tag{2.5}$$

Lemma2.4[14] If $J(t)$ is a non-increasing function on $[t_0, +\infty), t_0 \geq 0$ such that

$$J'(t)^2 \geq a + bJ(t)^{\frac{2+\frac{1}{\delta}}{\delta}}, \quad \forall t_0 \geq 0, \tag{2.6}$$

where $a > 0, b \in R$. Then there exists a finite time T^* such that

$$\lim_{t \rightarrow T^*} J(t) = 0 \tag{2.7}$$

And the case that

$$(i) \quad \text{If } b > 0, \quad J(t_0) < \min\left\{1, \sqrt{\frac{a}{-b}}\right\}, \text{ then}$$

$$T^* \leq t_0 + \frac{1}{\sqrt{-b}} \ln \frac{\sqrt{\frac{a}{-b}}}{\sqrt{\frac{a}{-b}} - J(t_0)}; \tag{2.8}$$

$$(ii) \quad \text{If } b = 0, \text{ then } T^* \leq t_0 + \frac{J(t_0)}{\sqrt{a}}. \tag{2.9}$$

$$(iii) \quad \text{If } b > 0, \quad T^* \leq \frac{J(t_0)}{\sqrt{a}} \quad \text{or}$$

$$T^* \leq t_0 + 2^{\frac{3\delta+1}{2\delta}} \frac{\delta c}{\sqrt{a}} \{1 - [1 + cJ(t_0)]^{\frac{-1}{2\delta}}\} \tag{2.10}$$

III. EXISTENCE OF LOCAL SOLUTION

In this section ,we prove the existence of local solution to problem (1.1)-(1.3) for initial value $(u_0, u_1) \in H^m(\Omega) \cap H^{2m}(\Omega) \times H^m(\Omega)$

Theorem 3.1 Suppose that (G₁), (G₂) hold, for any give $(u_0, u_1) \in H^{2m}(\Omega) \cap H_0^m(\Omega) \times L^2(\Omega)$

$$u(t) \in C([0, T_0]; H_0^m(\Omega) \cap H^{2m}(\Omega))$$

$$u_t(t) \in C([0, T_0]; L^2(\Omega) \cap H^m(\Omega))$$

Proof we set for any $T > 0$, the Banach space

$$X_{T,R} = \left\{ \begin{array}{l} v(t) \in C([0, T]; H^{2m} \cap H_0^m), \\ v_t(t) \in C([0, T]; L^2 \cap H_0^m), \\ \rho(v(t)) \leq R. \end{array} \right\}$$

Where $\rho(v(t)) = \|v_t\|^2 + \|(-\Delta)^m v\|^2$. Define distance $\rho(u, v) = \sup_{t \leq 0 \leq T} \rho(u(t) - v(t))$. Then $X_{T,R}$ is a complete metric space.

We define map the $\mathcal{X} : v \mapsto \mathcal{X}(v) = u$. u is the unique solution of the following equation:

$$u_{tt} + (-\Delta)^m u_t + \phi(\|\nabla^m u\|^2)(-\Delta)^m u + h(u_t) = f(v) \quad (3.1)$$

$$u(x, t) = 0, \quad \frac{\partial^i u(x, t)}{\partial v^i} = 0, \quad i = 1, 2, \dots, m-1, \quad x \in \partial\Omega, t \geq 0. \quad (3.2)$$

$$u(x, 0) = u_0(x), \quad u_t(x, 0) = u_1(x), \quad x \in \Omega. \quad (3.3)$$

we define map the $T > 0$ and $R > 0$, such that

- (1) \mathcal{X} maps $X_{T,R}$ into itself;
- (2) \mathcal{X} is a contraction mapping with respect to the metric $d(\dots)$.

Step 1. we will show that \mathcal{X} maps $X_{T,R}$ into itself, we multiply $u_t + \varepsilon(-\Delta)^m u$ with both sides of equation (3.4). and obtain

$$(u_{tt} + (-\Delta)^m u_t + \phi(\|\nabla^m u\|^2)(-\Delta)^m u + h(u_t), u_t + (-\Delta)^m u) = (f(v), u_t + (-\Delta)^m u) \quad (3.4)$$

$$\begin{aligned} & (u_{tt}, u_t + \varepsilon(-\Delta)^m u) \\ &= \frac{1}{2} \frac{d}{dt} \|u_t\|^2 + \frac{\varepsilon d}{dt} (u_t, (-\Delta)^m u) - \varepsilon \|\nabla^m u_t\|^2 \end{aligned} \quad (3.5)$$

$$\begin{aligned} & ((-\Delta)^m u_t, u_t + \varepsilon(-\Delta)^m u) \\ &= \|\nabla^m u_t\|^2 + \frac{\varepsilon}{2} \frac{d}{dt} \|(-\Delta)^m u\|^2 \end{aligned} \quad (3.6)$$

$$\begin{aligned} & (\phi(\|\nabla^m v\|^2)(-\Delta)^m u, u_t + \varepsilon(-\Delta)^m u) \\ &= \frac{1}{2} \frac{d}{dt} (\phi(\|\nabla^m v\|^2) \|\nabla^m u\|^2) \\ & - \frac{1}{2} \frac{d}{dt} (\phi(\|\nabla^m v\|^2) \|\nabla^m u\|^2 + \varepsilon \phi(\|\nabla^m v\|^2) \|(-\Delta)^m u\|^2) \end{aligned} \quad (3.7)$$

$$\begin{aligned} & \left| \frac{d}{dt} (\phi(\|\nabla^m v\|^2) \|\nabla^m u\|^2) \right| \\ & \leq \|\phi'(\xi)\|_\infty \|(-\Delta)^m v\| \|v_t\| \|\nabla^m u\|^2 \\ & \leq LR \|\nabla^m u\|^2 \\ & \leq \frac{LR}{\lambda^m} \|(-\Delta)^m u\|^2 \\ & \leq \kappa_0 \rho(u(t)) \end{aligned} \quad (3.8)$$

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$$\begin{aligned}
 & (\phi(\|\nabla^m v\|^2))(-\Delta)^m u, u_t + \varepsilon(-\Delta)^m u \\
 & \geq \frac{1}{2} \frac{d}{dt} (\phi(\|\nabla^m v\|^2) \|\nabla^m u\|^2) \\
 & - \kappa_0 \rho(u(t)) + \varepsilon \phi(\|\nabla^m v\|^2) \|(-\Delta)^m u\|^2
 \end{aligned} \tag{3.9}$$

$$\begin{aligned}
 & (h(u_t), u_t + \varepsilon(-\Delta)^m u) \\
 & = (h(u_t), u_t) + \varepsilon(h(u_t), (-\Delta)^m u)
 \end{aligned} \tag{3.10}$$

According to young inequality and G_2 , such that

$$(h(u_t), \varepsilon(-\Delta)^m u) \geq -\frac{\varepsilon \|h(u_t)\|^2}{2} - \frac{\varepsilon \|(-\Delta)^m u\|^2}{2} \tag{3.11}$$

$$\begin{aligned}
 & \|h(u_t)\|^2 \\
 & \leq C_0(1 + \|\nabla^m u_t\|)^2 \\
 & \leq 2C_0 + 2C_0 \|\nabla^m u_t\|^2
 \end{aligned} \tag{3.12}$$

$$\begin{aligned}
 & (h(u_t), (-\Delta)^m u) \\
 & \geq -C_0 - C_0 \varepsilon \|\nabla^m u_t\|^2 - \frac{\varepsilon}{2} \|(-\Delta)^m u\|^2
 \end{aligned} \tag{3.13}$$

According to Lemma 2.1 and (G_1) , such that

$$\begin{aligned}
 & (f(v), u_t + \varepsilon(-\Delta)^m) \\
 & \leq \int_{\Omega} (k_1 |v| + k_2 |v|^{r+1}) u_t dx + \varepsilon \int_{\Omega} (k_1 |v| + k_2 |v|^{r+1}) (-\Delta)^m u dx \\
 & \leq (k_1 \|v\| + k_2 \|v\|_{2r+2}^{r+1}) \|u_t\| + \varepsilon (k_1 \|v\| + k_2 \|v\|_{2r+2}^{r+1}) \|(-\Delta)^m u\| \\
 & \leq \left(\frac{k_1 \sqrt{R}}{\lambda^{2m}} + k_2 \kappa R^{r+1} \right) \|u_t\| + \left(\frac{\varepsilon k_1 \sqrt{R}}{\lambda^{2m}} + \varepsilon k_2 \kappa R^{r+1} \right) \|(-\Delta)^m u\|
 \end{aligned} \tag{3.14}$$

We take $\alpha_0 = \left\{ \frac{k_1 \sqrt{R}}{\lambda^{2m}} + k_2 \kappa R^{r+1}, \frac{\varepsilon k_1 \sqrt{R}}{\lambda^{2m}} + \varepsilon k_2 \kappa R^{r+1} \right\}$, such that

$$\begin{aligned}
 & (f(v), u_t + \varepsilon(-\Delta)^m) \\
 & \leq \alpha_0 (\|u_t\| + \|(-\Delta)^m u\|) \\
 & \leq \alpha_0 (\rho(u))^{\frac{1}{2}}
 \end{aligned} \tag{3.15}$$

From (3.5) and (3.15), we have

$$\begin{aligned}
 & \frac{1}{2} \frac{d}{dt} (\|u_t\|^2 + 2\varepsilon(u_t, (-\Delta)^m u) + \varepsilon \|(-\Delta)^m u\|^2 + \phi(\|\nabla^m v\|^2) \|\nabla^m u\|^2) \\
 & + \|\nabla^m u_t\|^2 - \varepsilon \|\nabla^m u_t\|^2 - C_0 \varepsilon \|\nabla^m u_t\|^2 - \frac{\varepsilon}{2} \|(-\Delta)^m u\|^2 \\
 & \leq \alpha_0 (\rho(u))^{\frac{1}{2}} + \kappa_0 \rho(u(t)) + C_0
 \end{aligned} \tag{3.16}$$

$$\begin{aligned}
 & \frac{dt}{d} \left(\frac{1}{2} \|u_t\|^2 + \varepsilon(u_t, (-\Delta)^m u) + \|(-\Delta)^m u\|^2 + \frac{1}{2} \phi(\|\nabla^m v\|^2) \|\nabla^m u\|^2 \right) + K_1 \\
 & + \left(1 - \varepsilon - \frac{C_0 \varepsilon}{2} \right) \|\nabla^m u_t\|^2 + \varepsilon \phi(\|\nabla^m v\|^2) \|(-\Delta)^m u\|^2
 \end{aligned}$$

$$\leq \alpha_0(\rho(u))^{\frac{1}{2}} + \kappa_0\rho(u(t)) + \frac{\varepsilon}{2}\rho(u(t)) + C_0 \quad (3.17)$$

We take $\rho_1(u(t)) = \|u_t\|^2 + 2\varepsilon(u_t, (-\Delta)^m u) + \varepsilon\|(-\Delta)^m u\|^2 + \phi(\|\nabla^m v\|^2)\|\nabla^m u\|^2 + 2K_1$ such that

$$\begin{aligned} & \|u_t\|^2 + 2\varepsilon(u_t, (-\Delta)^m u) + \varepsilon\|(-\Delta)^m u\|^2 + \phi(\|\nabla^m v\|^2)\|\nabla^m u\|^2 + 2K_1 \\ & \geq \|u_t\|^2 + 2\varepsilon(u_t, (-\Delta)^m u) + \varepsilon\|(-\Delta)^m u\|^2 \end{aligned} \quad (3.18)$$

$$2\varepsilon(u_t, (-\Delta)^m u) \geq 2\varepsilon\|u_t\|^2 - \frac{\varepsilon}{2}\|(-\Delta)^m u\|^2 \quad (3.19)$$

From (3.18) and (3.19), we have

$$\|u_t\|^2 + 2\varepsilon(u_t, (-\Delta)^m u) + \varepsilon\|(-\Delta)^m u\|^2 + \phi(\|\nabla^m v\|^2)\|\nabla^m u\|^2 + 2K_1 \quad (3.20)$$

$$\geq (1 - 2\varepsilon)\|u_t\|^2 + \frac{\varepsilon}{2}\|(-\Delta)^m u\|^2 \quad (3.21)$$

We take $\alpha_1 = \min\{1 - \varepsilon, \frac{\varepsilon}{2}\}$, such that

$$\rho_1(u(t)) \geq \alpha_1(\|u_t\|^2 + \|(-\Delta)^m u\|^2) \quad (3.22)$$

So

$$\rho_1(u(t)) \geq \alpha_1\rho(u(t)) \quad (3.23)$$

$$\begin{aligned} & \frac{dt}{d}(\|u_t\|^2 + 2\varepsilon(u_t, (-\Delta)^m u) + \varepsilon\|(-\Delta)^m u\|^2 + \phi(\|\nabla^m v\|^2)\|\nabla^m u\|^2 + 2K_1) + \alpha_1\|\nabla^m u_t\|^2 \\ & \leq \alpha_0(\rho(u))^{\frac{1}{2}} + \kappa_0\rho(u(t)) + \frac{\varepsilon}{2}\rho(u(t)) + C_0 \end{aligned} \quad (3.24)$$

Where $\alpha_2 = 1 - \varepsilon - \frac{\varepsilon C_0}{2}$, $C_0 \leq \frac{2\kappa_0 K_1}{\alpha_1}$

$$\begin{aligned} & \frac{d}{dt}\rho_1(u(t)) + \alpha_2\|\nabla^m u_t\|^2 \\ & \leq \frac{2\kappa_0}{\alpha_1}\rho_1(u(t)) + \frac{\varepsilon}{2}\rho_1(u(t)) + \alpha_0(\rho_1(u))^{\frac{1}{2}}. \end{aligned} \quad (3.25)$$

So, we using Gronwall inequality, we obtain

$$\begin{aligned} & \rho_1(u(t)) + \alpha_2 \int_0^t \|\nabla^m u_s\|^2 ds \\ & \leq \{\rho_1(u(0))^{\frac{1}{2}} + \alpha_0 T\}^2 \rho^{\frac{(4\kappa_0 + \alpha_0 \varepsilon)T}{2\alpha_1}}. \end{aligned} \quad (3.26)$$

Where $\rho_1(u(0)) = \|u_1\|^2 + 2\varepsilon(u_1, (-\Delta)^m u_0) + \varepsilon\|(-\Delta)^m u_0\|^2 + \phi(\|\nabla^m v_0\|^2)\|\nabla^m u_0\|^2 + 2K_1$

From we obtain

$$\begin{aligned} & \rho(u(t)) + \alpha_2 \int_0^t \|\nabla^m u_s\|^2 ds \\ & \leq \frac{1}{\alpha_1}\rho_1(u(t)) + \frac{\alpha_2}{\alpha_1} \int_0^t \|\nabla^m u_s\|^2 ds \end{aligned}$$

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$$\leq \{\rho_1(u(0))^2 + \alpha_0 T\}^2 \rho^{\frac{(4\kappa_0 + \alpha_0 \varepsilon)T}{2\alpha_1}}. \quad (3.27)$$

Therefore, we choose that parameters T , and R , we obtain

$$\{\rho_1(u(0))^2 + \alpha_0 T\}^2 \rho^{\frac{(4\kappa_0 + \alpha_0 \varepsilon)T}{2\alpha_1}} \leq R^2. \quad (3.28)$$

So, we finishes the proof of the first step.

Step 2 we will show that \mathcal{X} is a contraction in $X_{T,R}$, let $v_1, v_2 \in X_{T,R}$, such that $\mathcal{X}(v_1) = u_1$
 $\mathcal{X}(v_2) = u_2$, Setting $w = u_1 - u_2$, then we have

$$\begin{aligned} w_{tt} + (-\Delta)^m w_t + \phi(\|\nabla^m v_1\|^2)(-\Delta^m)u_1 - \phi(\|\nabla^m v_2\|^2)(-\Delta^m)u_2 + h(u_{t1}) - h(u_{t2}) &= f(v_1) - f(v_2) \\ (w_{tt}, w_t + (-\Delta)^m w) & \\ = \frac{1}{2} \frac{d}{dt} \|w_t\|^2 + \frac{\varepsilon d}{dt} (w_t, (-\Delta)^m w) - \varepsilon \|\nabla^m w_t\|^2 & \end{aligned} \quad (3.29)$$

$$\begin{aligned} ((-\Delta)^m w_t, w_t + \varepsilon(-\Delta)^m w) & \\ = \|\nabla^m w_t\|^2 + \frac{\varepsilon}{2} \frac{d}{dt} \|(-\Delta)^m w\|^2 & \end{aligned} \quad (3.30)$$

$$\begin{aligned} & (\phi(\|\nabla^m v_1\|^2)(-\Delta^m)u_1 - \phi(\|\nabla^m v_2\|^2)(-\Delta^m)u_2, w_t + \varepsilon(-\Delta)^m w) \\ &= (\phi(\|v_1\|^2)(-\Delta)^m w, w_t + \varepsilon(-\Delta)^m w) \\ &+ (\phi(\|\nabla^m v_1\|^2)(-\Delta)^m u_2 - \phi(\|\nabla^m v_2\|^2)(-\Delta)^m u_2, w_t + \varepsilon(-\Delta)^m w) \\ &= \frac{1}{2} \frac{d}{dt} \phi(\|\nabla^m v_1\|^2) \|\nabla^m w\|^2 + \phi(\|\nabla^m v_1\|^2) \|(-\Delta)^m w\|^2 - \frac{1}{2} \frac{d}{dt} (\phi(\|\nabla^m v_1\|^2)) \|\nabla^m w\|^2 \\ &+ (\phi'(\xi)(\|\nabla^m v_1\|^2 - \|\nabla^m v_2\|^2))(-\Delta)^m u_2, w_t) \\ &+ (\phi'(\xi)(\|\nabla^m v_1\|^2 - \|\nabla^m v_2\|^2))(-\Delta)^m u_2, \varepsilon(-\Delta)^m w) \\ &\geq \frac{1}{2} \frac{d}{dt} \phi(\|\nabla^m v_1\|^2) \|\nabla^m w\|^2 + \varepsilon \phi(\|\nabla^m v_1\|^2) \|(-\Delta)^m w\|^2 \\ &- \phi'(\|\nabla^m v_1\|^2) \int_{\Omega} \nabla^m v_1 \nabla^m v_{1t} dx \|\nabla^m w\|^2 + \|\phi'(\xi)\|_{\infty} (\|\nabla^m v_1\|^2 - \|\nabla^m v_2\|^2) (-\Delta)^m u_2, w_t) \\ &- \varepsilon \|\phi'(\xi)\|_{\infty} (\|\nabla^m v_1\| + \|\nabla^m v_2\|) (\|\nabla^m v_1\| - \|\nabla^m v_2\|) \|(-\Delta)^m u_2\| \|(-\Delta)^m w\| \\ &\geq \frac{1}{2} \frac{d}{dt} \phi(\|\nabla^m v_1\|^2) \|\nabla^m w\|^2 + \varepsilon \phi(\|\nabla^m v_1\|^2) \|(-\Delta)^m w\|^2 - L_0 \|(-\Delta)^m w\|^2 \\ &- \|\phi'(\xi)\|_{\infty} (\|\nabla^m v_1\| + \|\nabla^m v_2\|) (\|\nabla^m v_1\| - \|\nabla^m v_2\|) \|(-\Delta)^m u_2\| \|w_t\| \\ &- \varepsilon \|\phi'(\xi)\|_{\infty} (\|\nabla^m v_1\| + \|\nabla^m v_2\|) (\|\nabla^m v_1\| - \|\nabla^m v_2\|) \|(-\Delta)^m u_2\| \|(-\Delta)^m w\| \\ &\geq \frac{1}{2} \frac{d}{dt} \phi(\|\nabla^m v_1\|^2) \|\nabla^m w\|^2 + \varepsilon \phi(\|\nabla^m v_1\|^2) \|(-\Delta)^m w\|^2 - L_0 \|(-\Delta)^m w\|^2 \\ &- L_1 \|(-\Delta)^m (v_1 - v_2)\| \|(-\Delta)^m u_2\| \|w_t\| - C \|(-\Delta)^m w\|^2 \\ &- L_2 \|(-\Delta)^m (v_1 - v_2)\| \|(-\Delta)^m u_2\| \|(-\Delta)^m w\| \\ &\geq \frac{1}{2} \frac{d}{dt} \phi(\|\nabla^m v_1\|^2) \|\nabla^m w\|^2 + \varepsilon \phi(\|\nabla^m v_1\|^2) \|(-\Delta)^m w\|^2 \end{aligned}$$

$$-L_3[\rho(v_1 - v_2)]^{\frac{1}{2}}[\rho(w)]^{\frac{1}{2}} - L_0\rho(w(t)) \tag{3.31}$$

From G_2 we have

$$\begin{aligned} & (h(u_{t_1}) - h(u_{t_2}), w_t + \varepsilon(-\Delta)^m w) \\ & \leq L_4\|w_t\|^2 + L_5\|(-\Delta)^m w\|^2 \\ & \leq \beta_0\rho(w(t)) \end{aligned} \tag{3.32}$$

Where $\beta_0 = \max\{L_4, L_5\}$.

From G_1 and Lemma 2.1 we have

$$\begin{aligned} & (f(v_1) - f(v_2), w_t + \varepsilon(-\Delta)^m w) \\ & \leq \int_{\Omega} k_1(|v_1| - |v_2|)|w_t| + k_2(|v_1|^{r+1} - |v_2|^{r+1})|w_t| dx \\ & + \int_{\Omega} k_1(|v_1| - |v_2|)|w_t| + k_2(|v_1|^{r+1} - |v_2|^{r+1})|(-\Delta)^m w| dx \\ & \leq (k_1 + k_2(\|v_1\|_{nr}^r + \|v_2\|_{nr}^r))\|v_1 - v_2\|_{\frac{2n}{n-2}}\|(-\Delta)^m w\| \\ & + (k_1 + k_2(\|v_1\|_{nr}^r + \|v_2\|_{nr}^r))\|v_1 - v_2\|_{\frac{2n}{n-2}}\|w_t\| \\ & \leq L_6\|(-\Delta)^m(v_1 - v_2)\|\|w_t\| + L_7\|(-\Delta)^m(v_1 - v_2)\|\|(-\Delta)^m w\| \\ & \leq L_8[\rho(v_1 - v_2)]^{\frac{1}{2}}[\rho(w)]^{\frac{1}{2}} + L_9[\rho(v_1 - v_2)]^{\frac{1}{2}}[\rho(w(t))]^{\frac{1}{2}} \end{aligned} \tag{3.33}$$

From (3.29)-(3.33), we obtain

$$\begin{aligned} & \frac{d}{dt}\left(\frac{1}{2}\|w_t\|^2 + \frac{1}{2}\phi(\|\nabla^m v_1\|^2)\|\nabla^m w\|^2 + \varepsilon\frac{1}{2}\|(-\Delta)^m w\|^2 + \varepsilon(w_t, (-\Delta)^m w)\right) \\ & + \|\nabla^m w_t\|^2 - \varepsilon\|\nabla^m w_t\|^2 + \varepsilon\phi(\|\nabla^m v_1\|^2)\|(-\Delta)^m w\|^2 \\ & - L_3[\rho(v_1 - v_2)]^{\frac{1}{2}}[\rho(w)]^{\frac{1}{2}} - L_0\rho(w(t)) - \beta_0\rho(w(t)) \\ & \leq L_8[\rho(v_1 - v_2)]^{\frac{1}{2}}[\rho(w)]^{\frac{1}{2}} + L_9[\rho(v_1 - v_2)]^{\frac{1}{2}}[\rho(w(t))]^{\frac{1}{2}} \end{aligned} \tag{3.34}$$

From (3.34) we obtain

$$\begin{aligned} & \frac{d}{dt}(\|w_t\|^2 + \phi(\|\nabla^m v_1\|^2)\|\nabla^m w\|^2 + \varepsilon\|(-\Delta)^m w\|^2 + 2\varepsilon(w_t, (-\Delta)^m w)) \\ & \leq (L_0 + \beta_0)\rho(w(t)) + (L_3 + L_8 + L_9)[\rho(v_1 - v_2)]^{\frac{1}{2}}[\rho(w)]^{\frac{1}{2}} \end{aligned} \tag{3.35}$$

Where we set

$$\rho_2(w(t)) = \|w_t\|^2 + \phi(\|\nabla^m v_1\|^2)\|\nabla^m w\|^2 + \varepsilon\|(-\Delta)^m w\|^2 + 2\varepsilon(w_t, (-\Delta)^m w) \tag{3.36}$$

$$2\varepsilon(w_t, (-\Delta)^m w) \geq 2\varepsilon\|w_t\|^2 - \frac{\varepsilon}{2}\|(-\Delta)^m w\|^2 \tag{3.37}$$

Then we have

$$\begin{aligned} & \rho_2(w(t)) \\ & = \|w_t\|^2 + \phi(\|\nabla^m v_1\|^2)\|\nabla^m w\|^2 + \varepsilon\|(-\Delta)^m w\|^2 + 2\varepsilon(w_t, (-\Delta)^m w) \\ & \geq \|w_t\|^2 + \varepsilon\|(-\Delta)^m w\|^2 + 2\varepsilon(w_t, (-\Delta)^m w) \end{aligned}$$

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$$\begin{aligned} &\geq (1-2\varepsilon)\|w_t\|^2 + \frac{\varepsilon}{2}\|(-\Delta)^m w\|^2 \\ &\geq \beta_1(\|w_t\|^2 + \|(-\Delta)^m w\|^2) \end{aligned} \tag{3.38}$$

Where $\beta_1 = \min\{1-2\varepsilon, \frac{\varepsilon}{2}\}$

$$\begin{aligned} &\frac{d\rho_2(w(t))}{dt} \\ &\leq (L_0 + \beta_0)\rho(w(t)) + (L_3 + L_8 + L_9)[\rho(v_1 - v_2)]^{\frac{1}{2}}[\rho(w)]^{\frac{1}{2}} \end{aligned} \tag{3.39}$$

$$\leq \left(\frac{L_0 + \beta_0}{\beta_1}\right)\rho_2(w(t)) + (L_3 + L_8 + L_9) \frac{1}{\beta_1} \{\rho_2(v_1(t) - v_2(t))\}^{\frac{1}{2}} \{\rho_2(w(t))\}^{\frac{1}{2}} \tag{3.40}$$

Where $\rho_2(w(0)) = 0$.

Therefore, we use Gronwall inequality, we have

$$\rho_2(w(t)) \leq (L_3 + L_8 + L_9)^2 T^2 e^{\frac{(L_0 + \beta_0)T}{\beta_1}} \sup_{0 \leq t \leq T} \rho(v_1(t) - v_2(t)) \tag{3.41}$$

So, we have

$$\sup_{0 \leq t \leq T} \rho(u_1(t) - u_2(t)) \leq C_{T,R} \sup_{0 \leq t \leq T} \rho(v_1(t) - v_2(t)) \tag{3.42}$$

Where $C_{T,R} = (L_3 + L_8 + L_9)^2 T^2 e^{\frac{(L_0 + \beta_0)T}{\beta_1}}$.

It is easy to see $C_{T,R} < 1$ for a small $T > 0$.

we finish the proof of the second step. By applying the Banach contraction mapping theorem, the proof of Theorem 3.1 is now complete.

IV. THE BLOW-UP IN FINITE TIME

In this section, we consider the blow-up of solution, we assume that $\phi(s) := s^q$, $q \geq 0$, $h(u_t) = u_t$, such problems (1.1)-(1.3) became

$$u_{tt}(t) + (-\Delta)^m u_t(t) + \|\nabla^m u\|^{2q} (-\Delta)^m u(t) + u_t(t) = f(u(t)) \tag{4.1}$$

Next, we define the energy function of the solution u of (4.1)

$$E(t) = \frac{1}{2}\|u_t\|^2 + \frac{1}{2(q+1)}\|\nabla^m u\|^{2(q+1)} - \int_{\Omega} G(u)dx, \tag{4.2}$$

$$G(s) = \int_0^s g(s)ds.$$

where

Then, we have

$$E(t) = E(0) - \int_0^t (\|\nabla^m u_t\|^2 + \|u_t\|^2) ds, \tag{4.3}$$

$$E(0) = \frac{1}{2}\|u_1\|^2 + \frac{1}{2(q+1)}\|\nabla^m u_0\|^{2(q+1)} - \int_{\Omega} G(u_0)dx.$$

where

Definition 4.1 A solution $u(t)$ of (4.1) is called a blow-up solution if there exist a finite time T^* such that

$$\lim_{t \rightarrow T^*} \int_{\Omega} (\|\nabla^m u\|^2 + |u|^2) dx = +\infty. \tag{4.4}$$

For the next lemma, we define

$$Y(t) = Y(u(t)) = \|u(t)\|^2 + \int_0^t (\|u(s)\|^2 + \|\nabla^m u(s)\|^2) ds, \quad \text{for } t \geq 0. \quad (4.5)$$

Lemma 4.1 Assume (G_1) , (G_3) hold, Let $\frac{2+4\beta}{q+1} - 2 \geq 0$, then we have

$$Y''(t) - 4(\beta+1)\|u_t\|^2 \geq (-4-8\beta)E(0) + (4+8\beta) \left[\int_0^t (\|u_t(s)\|^2 + \|\nabla^m u_t(s)\|^2) ds \right]. \quad (4.6)$$

Proof. From (4.5), we have

$$Y'(t) = 2 \int_{\Omega} u u_t dx + \|u(t)\|^2 + \|\nabla^m u\|^2, \quad (4.7)$$

and

$$Y''(t) = 2\|u_t(t)\|^2 - 2\|\nabla^m u\|^{2q+2} + 2 \int_{\Omega} f(u) u dx. \quad (4.8)$$

From above equation and the energy identity, we have

$$\begin{aligned} Y'' - 4(\beta+1)\|u_t\|^2 &= (-4-8\beta)E(0) + (4+8\beta) \int_0^t (\|u_t\|^2 + \|\nabla^m u\|^2) ds \\ &+ \int_{\Omega} (2f(u)u - (4+8\beta)G(u)) dx + \left\{ \frac{2+4\beta}{(q+1)} \|\nabla^m u\|^{2(q+1)} - 2\|\nabla^m u\|^{2q+2} \right\}, \end{aligned} \quad (4.9)$$

there form the assumption (G_3) , we have

$$2 \int_{\Omega} (f(u)u - (2+4\beta)G(u)) dx + \left[\frac{2+4\beta}{q+1} \|\nabla^m u\|^{2(q+1)} - 2\|\nabla^m u\|^{2q+2} \right] \geq 0. \quad (4.10)$$

So, we obtain (4.6)

Now, we can consider there different cases on the sign of initial energy $E(0)$

If $E(0) < 0$, from (4.6), we have

$$Y''(t) \geq (-4-8\beta)E(0), \quad (4.11)$$

integration (4.11) over $[0, t]$, we have that

$$Y'(t) \geq Y'(0) - 4(1+2\beta)E(0)t. \quad (4.12)$$

Thus, we get $Y'(t) > \|u_0\|^2 + \|\nabla^m u_0\|^2$ for $t > t^*$, where

$$t^* = \max \left\{ \frac{Y'(0) - (\|u_0\|^2 + \|\nabla^m u_0\|^2)}{4(1+\beta)E(0)}, 0 \right\}. \quad (4.13)$$

If $E(0) = 0$, from (4.6) we have

$$Y'' \geq 4(\beta+1)\|u_t\|^2 + (4+\beta) \int_0^t (\|u_t\|^2 + \|\nabla^m u\|^2) ds \geq 0 \quad t \geq 0 \quad (4.14)$$

Integration (4.9) over $[0, 1]$, we have

$$Y' \geq Y'(0) \quad (4.15)$$

Furthermore, if $Y'(0) > \|u_0\|^2 + \|\nabla^m u_0\|^2$, i.e. $\int_{\Omega} u_0 u_1 dx > 0$, so

$$Y'(t) \geq Y'(0) > \|\nabla^m u_0\|^2 + \|u_0\|^2, \quad t \geq 0. \quad (4.16)$$

If $E(0) > 0$ we first note that

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$$2\left(\int_0^t \int_{\Omega} uu_t dx dt + \int_0^t \int_{\Omega} \nabla^m u \nabla^m u_t dx ds\right) = \|u\|^2 + \|\nabla^m u\|^2 - (\|u_0\|^2 + \|\nabla^m u_0\|^2). \tag{4.10}$$

By using Holder inequality and (4.10) we have

$$\|u(t)\|^2 + \|\nabla^m u(t)\|^2 \leq \|u_0\|^2 + \|\nabla^m u_0\|^2 + \int_0^t \|u(s)\|^2 ds + \int_0^t \|u_t(s)\|^2 ds + \int_0^t \|\nabla^m u(s)\|^2 ds + \int_0^t \|\nabla^m u_t(s)\|^2 ds$$

so, from above, we obtain

$$\begin{aligned} Y'(t) &= 2\int_{\Omega} uu_t dx + \|u(t)\|^2 + \|\nabla^m u\|^2 \leq \|u(t)\|^2 + \|u_t(t)\|^2 + \|u(t)\|^2 + \|\nabla^m u(t)\|^2 \\ &\leq \|u_0\|^2 + \|\nabla^m u_0\|^2 + \int_0^t \|u(s)\|^2 ds + \int_0^t \|u_t(s)\|^2 ds + \int_0^t \|\nabla^m u(s)\|^2 ds + \int_0^t \|\nabla^m u_t(s)\|^2 ds \\ &\quad + \|u(t)\|^2 + \|u_t(t)\|^2 \end{aligned} \tag{4.11}$$

$$Y'(t) \leq Y(t) + \|u_t(t)\|^2 + \|\nabla^m u_0\|^2 + \|u_0\|^2 + \int_0^t \|\nabla^m u_t(s)\|^2 ds + \int_0^t \|u_t(s)\|^2 ds \tag{4.13}$$

Form above inequality and (4.6), we obtain

$$Y''(t) - 4(\beta + 1)Y'(t) + 4(1 + \beta)Y(t) + Y_1 \geq 0. \tag{4.14}$$

Where

$$Y_1 = (4 + 8\beta)E(0) + 4(1 + \beta)(\|u_0\|^2 + \|\nabla^m u_0\|^2). \tag{4.18}$$

Setting

$$B(t) = Y(t) + \frac{Y_1}{4(1 + \beta)} \quad t > 0 \tag{4.19}$$

Then $B(t)$ satisfies Lemma 2.3, if

$$Y'(0) > r_2 \left[Y(0) + \frac{Y_1}{4(1 + \beta)} \right] + \|u_0\|^2 + \|\nabla^m u_0\|^2 \tag{4.20}$$

Then $Y'(0) > \|u_0\|^2 + \|\nabla^m u_0\|^2$ for all $t > 0$

Lemma 4.2 Assume that (G_1) and (G_3) hold and that either one of the following conditions is satisfied $E(0) < 0$;

$$E(0) = 0 \quad \text{and} \quad \int_{\Omega} u_0 u_1 dx > 0;$$

$$E(0) > 0 \quad \text{and} \quad (4.20) \text{ holds, then } Y'(0) > \|u_0\|^2 + \|\nabla^m u_0\|^2 \text{ for all } t > 0$$

where $t = t^*$ is given by (4.13) in case(1); $t_0 = 0$ in cases (2) and (3).

Proof. we can prove this Lemma 4.2 by Lemma 4.1.

Theorem 4.1 Under the assumptions (G_1) and (G_3) , and that either one of the following condition is satisfied: $E(0) < 0$;

$$E(0) = 0 \quad \text{and} \quad \int_{\Omega} u_0 u_1 dx > 0;$$

$$0 < E(0) < \frac{\left(\int_{\Omega} u_0 u_1 dx\right)^2}{2(T_1 + 1)(\|\nabla^m u\|^2 + \|u_0\|^2)}, \quad \text{and by (4.20) hold.}$$

where T_1 to be chosen later

Then the solution u blow-up at finite T^* , and T^* can be estimate by (4.35)-(4.38), respectively, according to the sign of $E(0)$.
 Proof. Let

$$J(t) = (Y(t) - (\|u_0\|^2 + \|\nabla^m u_0\|^2))^{-\beta} \quad t \in [0, T_1] \quad (4.21)$$

where T_1 is some certain constant which will be chosen later. Then we get

$$J'(t) = -\beta J(t)^{1+\frac{1}{\beta}} (Y'(t) + (T_1 - t)(\|u_0\|^2 + \|\nabla^m u_0\|^2)) \quad (4.22)$$

And

$$J''(t) = -\beta J(t)^{1+\frac{2}{\beta}} V(t) \quad (4.23)$$

$$V(t) = Y''(t)[Y(t) + (T_1 - t)(\|u_0\|^2 + \|\nabla^m u_0\|^2)] - (1 + \beta)(Y'(t) - (\|u_0\|^2 + \|\nabla^m u_0\|^2)) \quad (4.24)$$

Next we denote

$$P = \|u(t)\|^2, \quad Q = \int_0^t \|\nabla^m u(s)\|^2 ds, \quad R = \|u_t(t)\|^2, \quad S = \int_0^t \|\nabla^m u_t(s)\|^2 ds \quad (4.25)$$

By (4.7) and (4.9), and Holder inequality we have

$$\begin{aligned} Y'(t) &= 2 \int_{\Omega} u(t)u_t(t) dx + \|u_0\|^2 + \|\nabla^m u_0\|^2 + 2 \int_0^t \int_{\Omega} (u(s)u_t(s) + \nabla^m u(s)\nabla^m u_t(s)) dx ds \\ &\leq \|u_0\|^2 + \|\nabla^m u_0\|^2 + 2\|u(t)\|\|u_t(t)\| + 2(\lambda^{-2m} + 1) \left(\int_0^t \|\nabla^m u_t\|^2 ds \right)^{\frac{1}{2}} \left(\int_0^t \|\nabla^m u\|^2 ds \right)^{\frac{1}{2}} \\ &\leq \|u_0\|^2 + \|\nabla^m u_0\|^2 + 2\sqrt{PR} + 2(\lambda^{-2m} + 1)\sqrt{QS} \\ &\leq \|u_0\|^2 + \|\nabla^m u_0\|^2 + \gamma(\sqrt{PR} + \sqrt{QS}) \end{aligned} \quad (4.26)$$

where $\gamma = 2 + 2\lambda^{-2m}$

From (4.6), we have

$$Y''(t) \geq (-4 - 8\beta)E(0) + (4 + 8\beta)(R + S) \quad (4.27)$$

From (4.24)-(4.27), we have

$$\begin{aligned} V(t) &\geq [(-4 - 8\beta)E(0) + 4(1 + \beta)(R + S)](Y(t) + (T_1 - t)(\|u_0\|^2 + \|\nabla^m u_0\|^2)) \\ &\quad - (1 + \beta)(Y'(t) - (\|u_0\|^2 + \|\nabla^m u_0\|^2)) \\ &\geq (-4 - 8\beta)E(0)J(t)^{\frac{-1}{\beta}} + 4(1 + \beta)(R + S)(T_1 - t)(\|u_0\|^2 + \|\nabla^m u_0\|^2) \\ &\quad + 4(1 + \beta)[(R + S)(P + Q) - (\sqrt{PR} + \sqrt{QS})^2] \\ &\geq (-4 - 8\beta)E(0)J(t)^{\frac{-1}{\beta}}. \quad t \geq t_0, \end{aligned} \quad (4.28)$$

from (4.23) we obtain

$$J''(t) \leq \beta(4 + 8\beta)E(0)J(t)^{1+\frac{1}{\beta}}. \quad t \geq t_0 \quad (4.29)$$

Now that by Lemma 2.4 that $J'(t) < 0$ for $t > t_0$ multiplying (4.29) by $J'(t)$ and integrating it from $t_0 \rightarrow t$, we have

$$J'(t) \geq \omega + \kappa J(t)^{\frac{2+1}{\beta}}, \quad t \geq t_0, \quad (4.30)$$

Where

$$\omega = \beta^2 J(t_0)^{\frac{2+2}{\beta}} [Y'(t_0) - (\|u_0\|^2 + \|\nabla^m u_0\|^2) - 8E(0)J(t_0)^{\frac{-1}{\beta}}] \quad (4.31)$$

and $\kappa = 8\beta^2 E(0).$ (4.33)

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We observe that

$$E(0) < \frac{(Y'(t_0) - (\|u_0\|^2 + \|\nabla^m u_0\|^2))}{8[Y(t_0)] + (T_1 - t_0)(\|u_0\|^2 + \|\nabla^m u_0\|^2)}$$

$\omega > 0$ if and if only

$$\lim_{t \rightarrow T^*} J(t) = 0$$

Then by Lemma4.2 ,there exists a finite time T^* ,such that and the upper bounds of T^* are estimated ,respectively ,according to the sign of $E(0)$,we obtain

$$\lim_{t \rightarrow T^*} (\|u(t)\|^2 + \int_0^t (\|u(s)\|^2 + \|\nabla^m u(s)\|^2)) = +\infty \tag{4.34}$$

The upper bounds of T^* are estimate as follows by Lemma4.2
In case (1),we have

$$T^* \leq t_0 - \frac{J(t_0)}{J'(t_0)} \tag{4.35}$$

Furthermore if $J(t_0) < \min\{1, \sqrt{\frac{\omega}{-\kappa}}\}$,then we obtain

$$T^* \leq t_0 + \frac{1}{\sqrt{-\kappa}} \operatorname{Ln} \frac{\sqrt{\frac{\omega}{-\kappa}}}{\sqrt{\frac{\omega}{-\kappa}} - J(t_0)} \tag{4.36}$$

In case (2),

$$T^* \leq t_0 - \frac{J(t_0)}{J'(t_0)} \quad \text{or} \quad T^* \leq t_0 - \frac{J(t_0)}{\sqrt{\omega}} \tag{4.37}$$

In case (3)

$$T^* \leq \frac{J(t_0)}{\sqrt{\omega}} \quad \text{or} \quad T^* \leq t_0 + 2^{\frac{3\beta+1}{2\beta}} \frac{\beta c}{\sqrt{\omega}} \{1 - [1 + cJ(t_0)]^{\frac{-1}{2\beta}}\} \tag{4.38}$$

Where $c = \left(\frac{\kappa}{\omega}\right)^{\frac{\beta}{2+\beta}}$.note that in case(1), $t_0 = t^*$ is given in (4.13),and in case(2) and case(3) $t_0 = 0$

Remark4.1 that in we observe that the choice of T_1 in (4.21) is feasible under the same condition as in[15].

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